

# Runge-Kutta Methods for solving ODEs



# ODE Solvers

## Runge Kutta

- Euler Method (1st order)
- Midpoint Method (2nd order)
- RK4 (4th order) - SciPy library
- RK45 (adaptive time step 4th and 5th order) - SciPy library - **BEST general purpose solver**

## Symplectic

- Euler-Cromer-Aspel Method (1st order)
- Leapfrog (2nd order) - **BEST for Hamiltonian (including Energy-Conserving) systems**

## Stiff Solvers

- Radau (5th order) - SciPy library - **Best all-purpose solver for stiff equations**
- BDF (order 1-5) - SciPy library

# Adaptive Time Step Runge Kutta 45

# Adaptive Time Step - Runge Kutta 45

## Method:

1. Step forward using RK4 method  $y_{RK4,n+1}$
2. Recalculate using a 5th-order Runge Kutta method  $y_{RK5,n+1}$
3. Find error between the two methods:  $\epsilon = |y_{RK4,n+1} - y_{RK5,n+1}|$
4. Test if  $\epsilon > atol + rtol \cdot |y_n|$ , then reduce the time step  $\Delta t$  by half and go back to step 1.
  - $atol$  = absolute tolerance
  - $rtol$  = relative tolerance
5. Test if  $\epsilon \leq atol + rtol \cdot |y_n|$ , then accept the value  $y_{RK4,n+1}$  and go to the next step

# Use the SciPy library to implement the RK45 method

1. Load solve\_ivp from SciPy library

```
from scipy.integrate import solve_ivp
```

2. Call solve\_ivp to perform integration

```
sol = solve_ivp(deriv_sho, (0, tmax), y0, method='RK45', args=params,  
atol=1e-4, rtol=1e-3)
```

derivative function

time limits

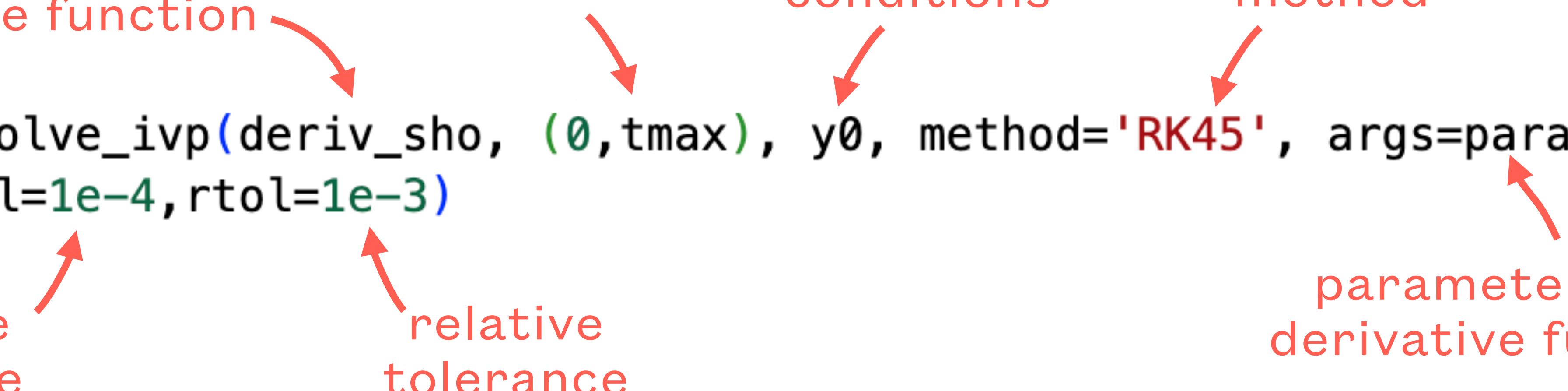
initial conditions

integration method

absolute tolerance

relative tolerance

parameters for derivative function



## Use the SciPy library to implement the RK45 method

```
sol = solve_ivp(deriv_sho, (0,tmax), y0, method='RK45', args=params,  
atol=1e-4,rtol=1e-3)
```

3. Extract time and solution from the solution object sol:

```
t = sol.t          # extract times  
x = sol.y[0,:]    # extract positions  
v = sol.y[1,:]    # extract velocities
```

# Complete code to implement RK45 integration

Solution

```
import numpy as np
import matplotlib.pyplot as plt
from scipy.integrate import solve_ivp

##### Parameters #####
m = 1          # mass
k = 1          # spring constant
tmax = 50      # maximum time
dt = 0.01      # time step
x0 = 1          # initial position
v0 = 0          # initial velocity

params = np.array([m,k])      # bundle derivative parameters together
y0 = np.array([x0,v0])        # bundle initial conditions together

##### Perform RK45 Integration #####
sol = solve_ivp(deriv_sho, (0,tmax), y0, method='RK45', args=params,
                atol=1e-4, rtol=1e-3)
t = sol.t          # extract times
x = sol.y[0,:]    # extract positions
v = sol.y[1,:]    # extract velocities

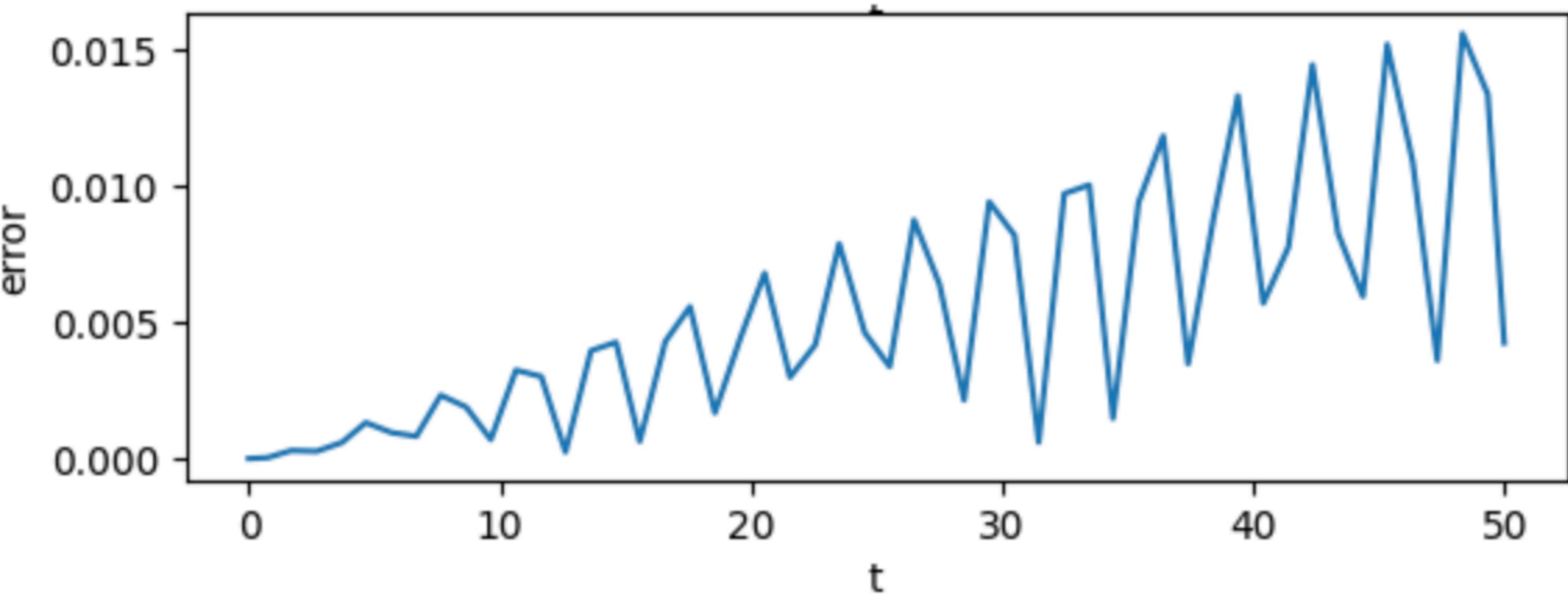
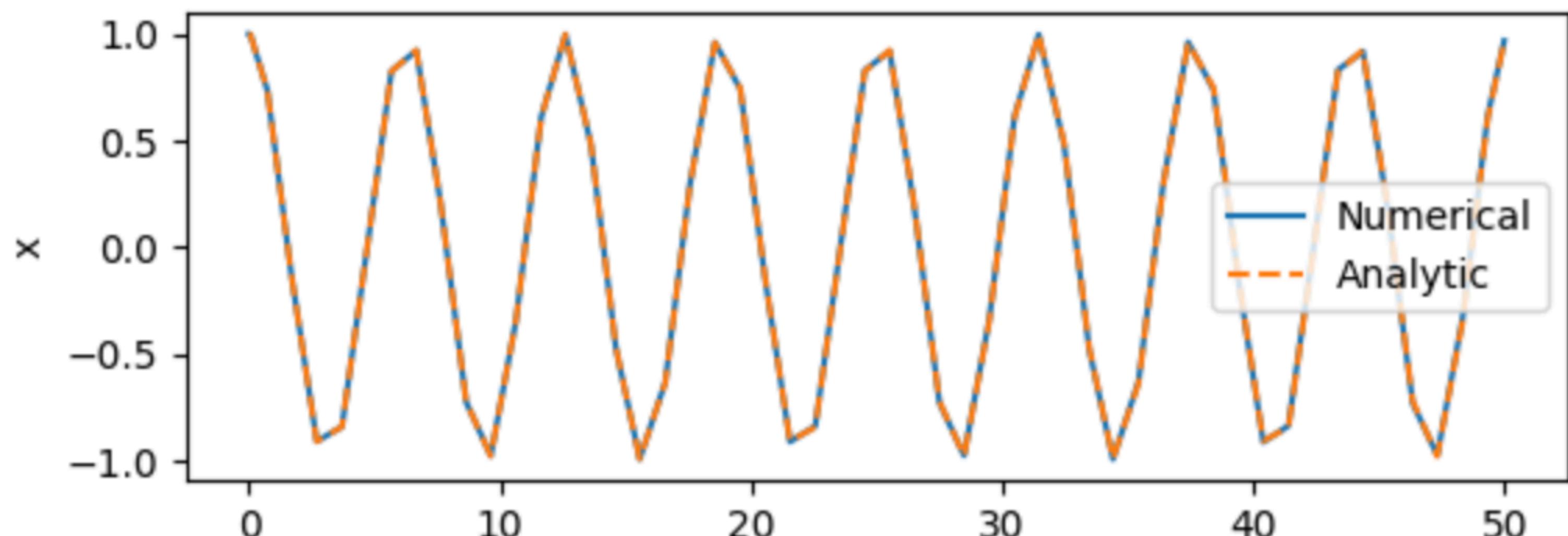
##### Analytic Solution #####
omega = np.sqrt(k/m)
x_true = x0 * np.cos(omega*t)

##### Plot Solution #####
plot_solution(x, x_true, t, "SH0 - RK45")
```

Error is set by the absolute and relative tolerances.

Times are unevenly distributed

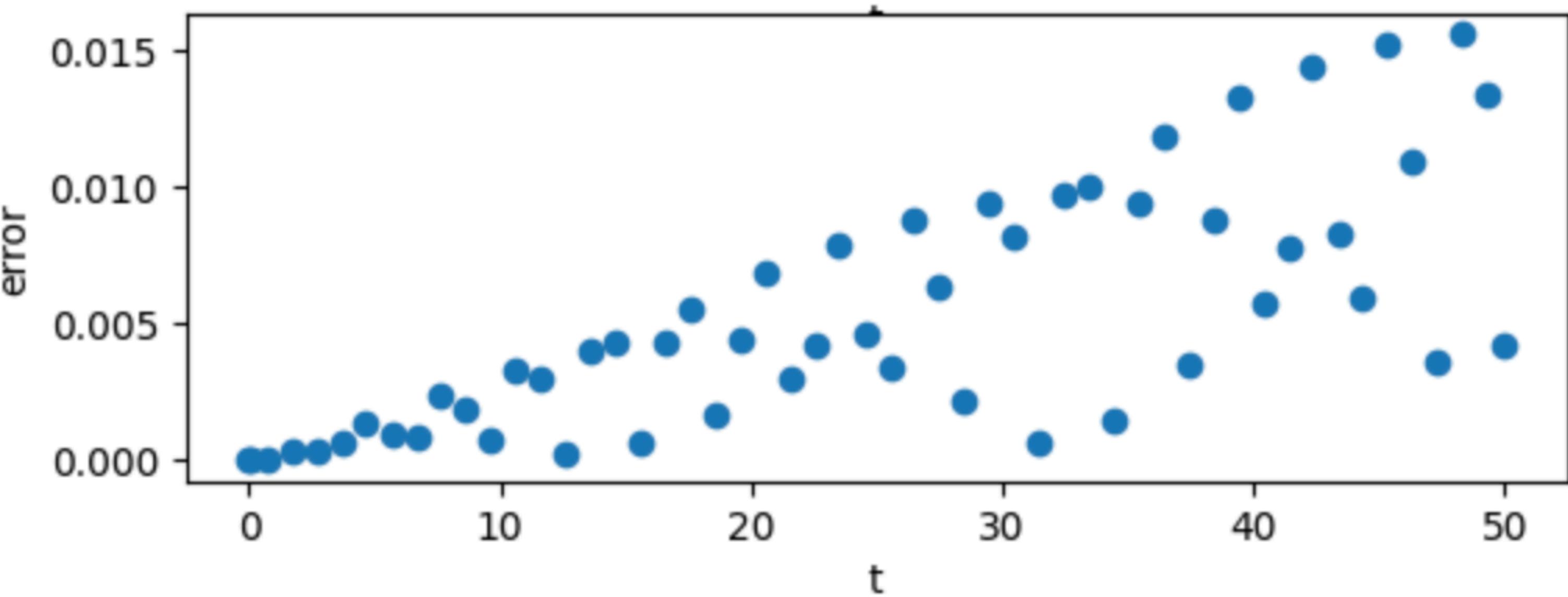
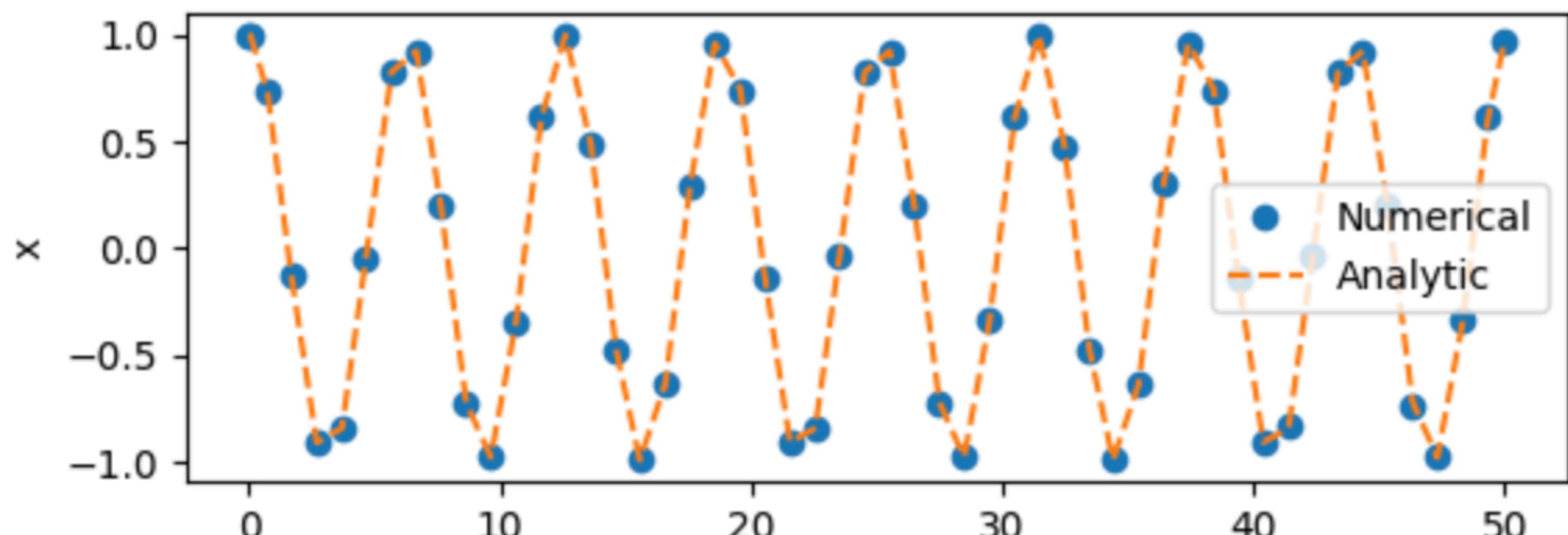
SHO - RK45



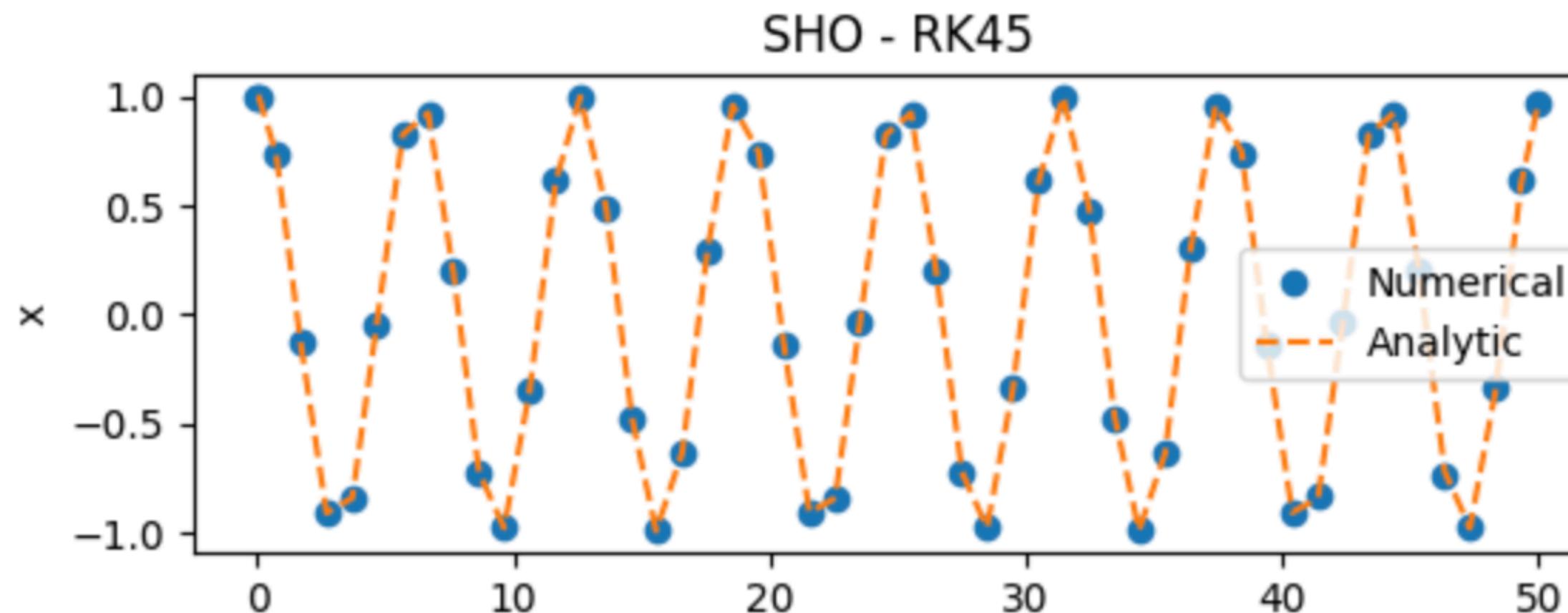
Error is set by the absolute and relative tolerances.

Times are unevenly distributed

SHO - RK45



# Did the RK45 do a good job ?



Two factors:

**Accuracy** - how well the numerical solution tracks the true solution

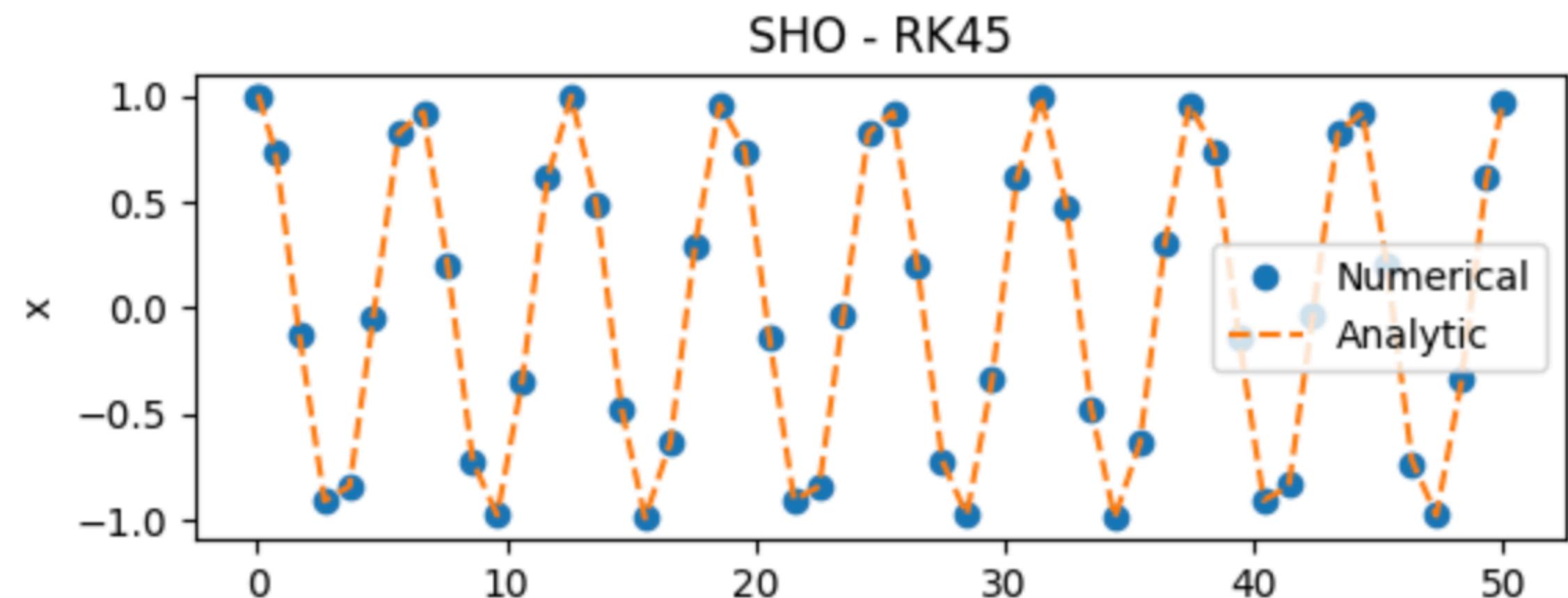
**Sampling** - how often we sample the solution effects the appearance of the graph.

- The error between the numerical solution and the analytic solution is around 1.5% of the oscillation's amplitude. This error could be reduced by imposing **stricter tolerances**.
- The plotted solution looks “jagged” because of poor sampling. We will show ways to use the numerical results to “fill out” the full curve using **interpolation**.

# Custom Evaluation Times

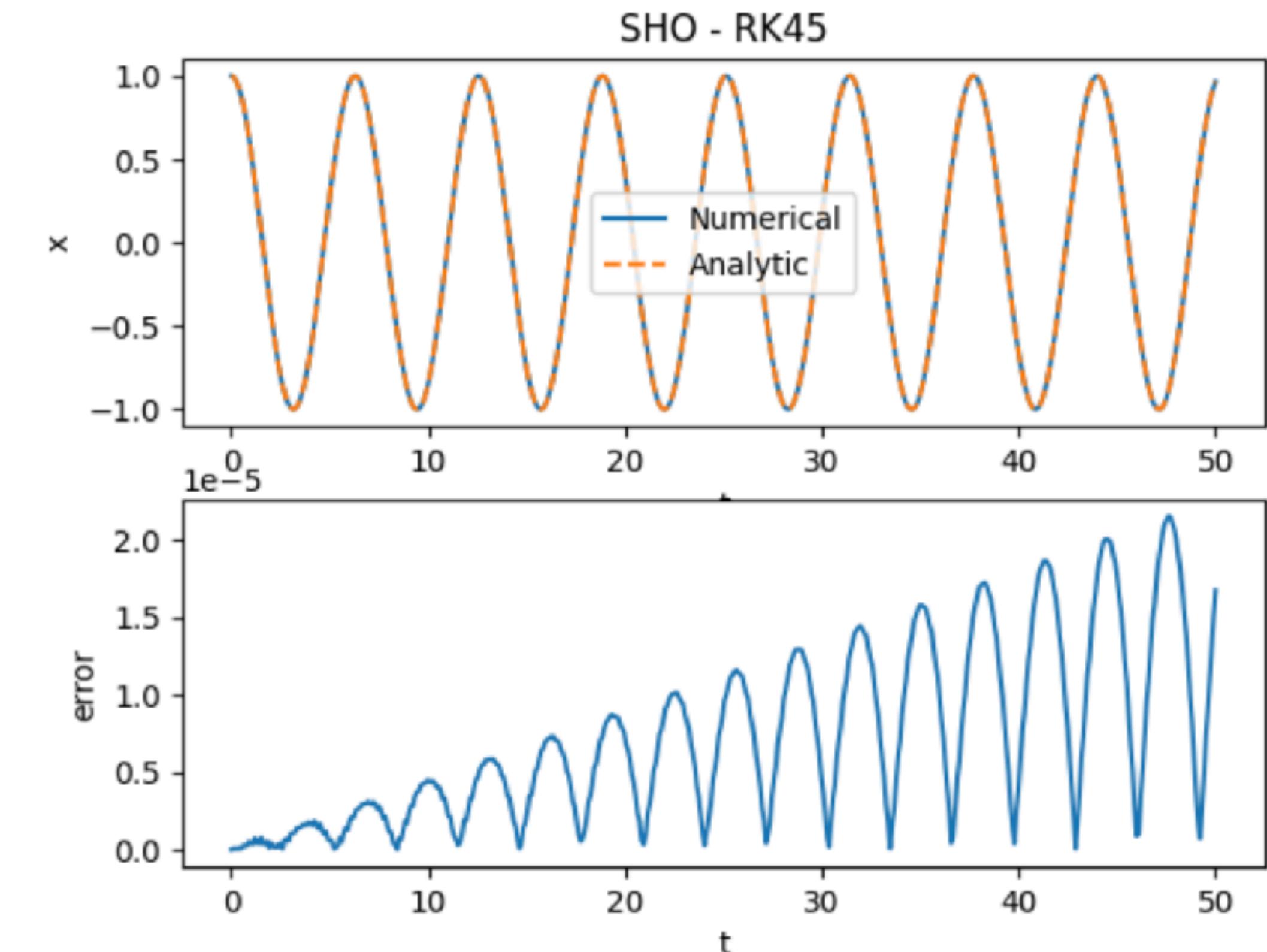
Adaptive time step methods produce solutions at nonuniform times. If you want to generate a solution on uniformly spaced times, you have two options:

- **option 1:** Pass an optional array of times to the `solve\_ivp()` function. This option modifies the time-stepping algorithm to include the specified times
- **option 2:** Turn on the `dense\_output` option. This option does not affect the numerical integration, but rather allows the user to use interpolation to evaluate the solution at any time they want.



# Custom Evaluation Times

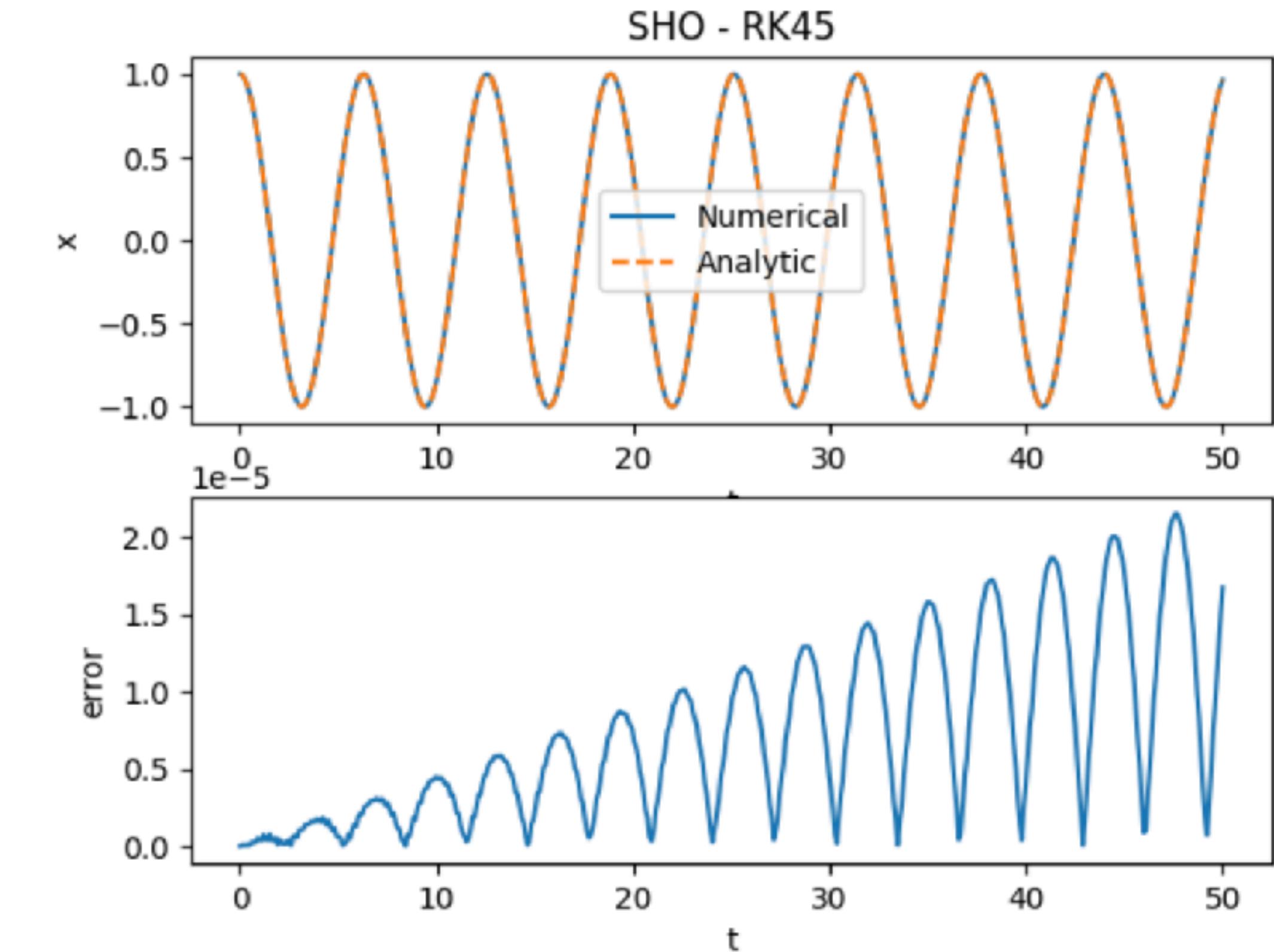
- **option 1:** Pass an optional array of times to the `solve\_ivp()` function. This option modifies the time-stepping algorithm to include the specified times



```
times = np.linspace(0,tmax,500)
sol = solve_ivp(deriv_sho, (0,tmax), y0, method='RK45', t_eval=times,
                args=params, atol=1e-6, rtol=1e-6)
t = sol.t          # extract times
x = sol.y[0,:]    # extract positions
v = sol.y[1,:]    # extract velocities
```

# Custom Evaluation Times

- **option 2:** Turn on the `dense\_output` option. This option does not affect the numerical integration, but rather allows the user to use interpolation to evaluate the solution at any time they want.

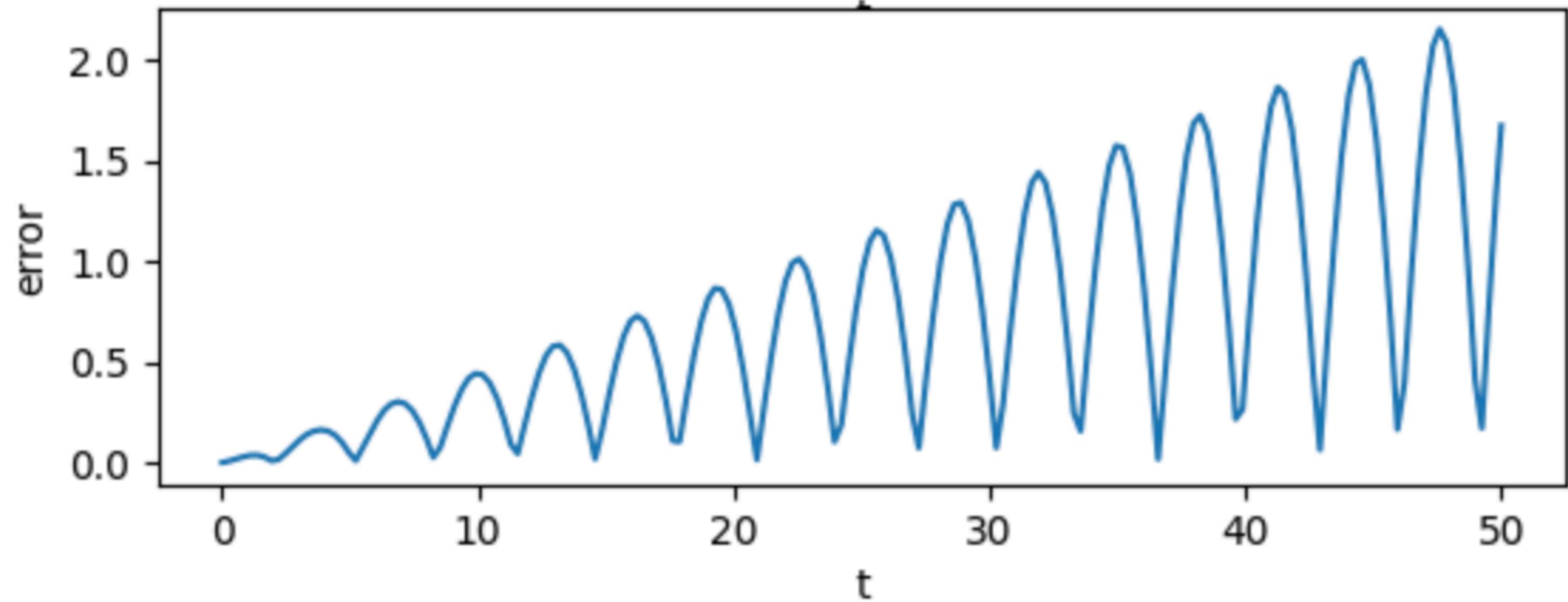
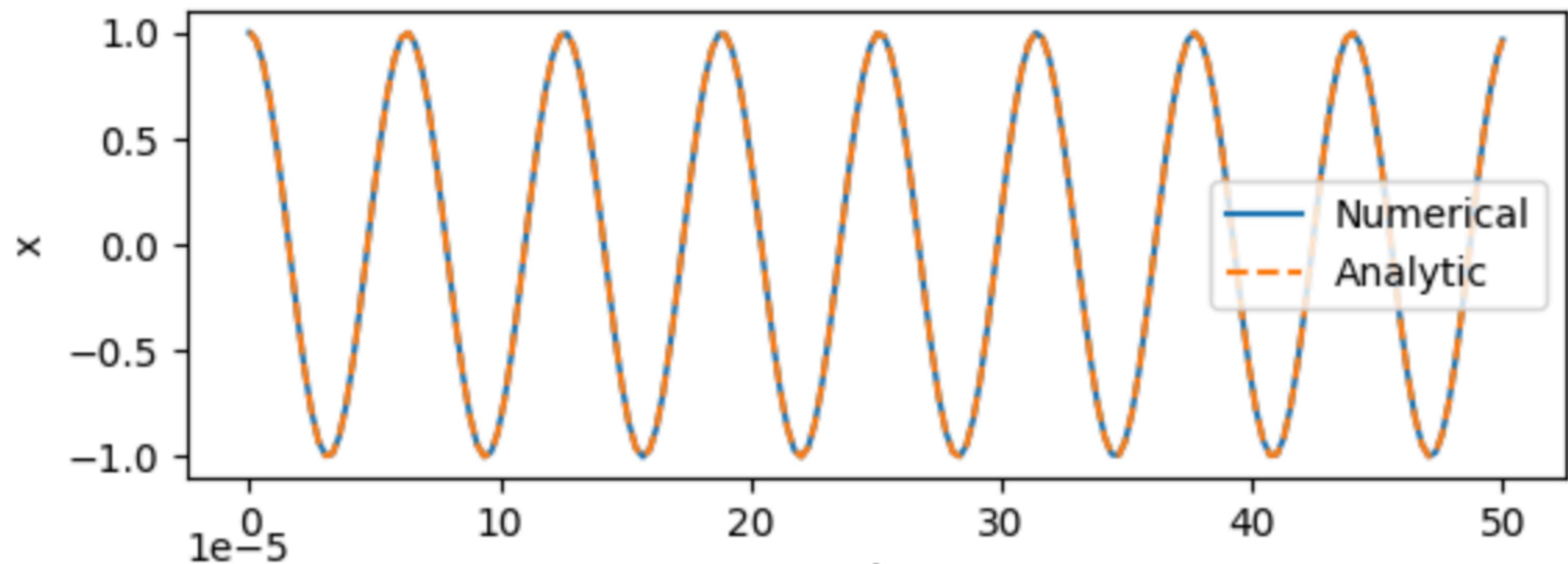


```
sol = solve_ivp(deriv_sho, (0,tmax), y0, method='RK45', dense_output=True,  
                args=params, atol=1e-6, rtol=1e-6)  
  
# Extract solution on a regular grid of time values  
t = np.linspace(0, tmax, 500)      # define time array  
y = sol.sol(t)                  # create a 2D solution array  
x = y[0,:]                      # extract position from the solution arr
```

Error is set by the absolute and relative tolerances.

Times are unevenly distributed

SHO - RK45



# Application: Chaos Theory



## Edward Lorenz

(1917-2008 )

- MIT Meteorologist
- Founder of Chaos Theory



## Margaret Hamilton

(1936- )

- Director MIT Instrument laboratory
- Developed software for NASA's Apollo Guidance Computer
- Coined term "software engineer"

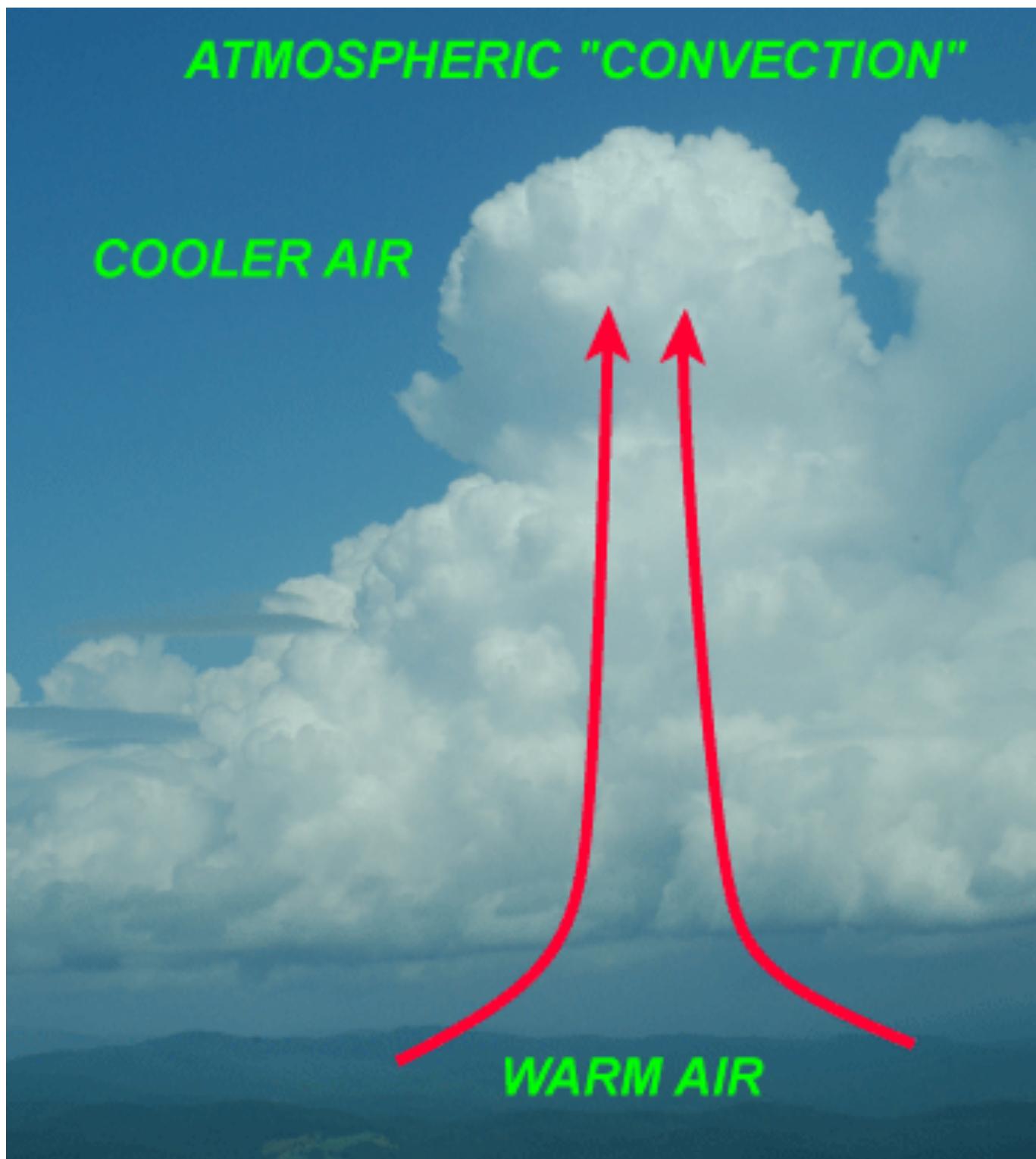


## Ellen Fetter

- Programmer
- Worked on Numerical Chaos Theory with Lorenz

# Lorenz System

Lorenz was looking for a simplified system of equations to model convection in the atmosphere, when he discovered this set of equations:



$$\frac{dx}{dt} = \sigma(y - x)$$

$x$  = flow rate of convection

$$\frac{dy}{dt} = x(r - z) - y$$

$y$  = Temp difference

$$\frac{dz}{dt} = xy - bz$$

$z$  = nonlinearity of Temp profile

$\sigma$  = Prandtl number

$r$  = Rayleigh number

$b$  = scale of the flow

Lorenz invented **Chaos Theory** to describe the dynamics of this system

# Dynamical Systems

Consider the dynamics of systems that have both:

- dissipation (energy removed from the system)
- external driving (energy added to system)

These systems often have “**attractors**” which are values toward which the system evolves over time. These attractors have different forms:

- fixed (equilibrium) points - 0D
- limit cycles - 1D
- strange attractors (fractal dimension)

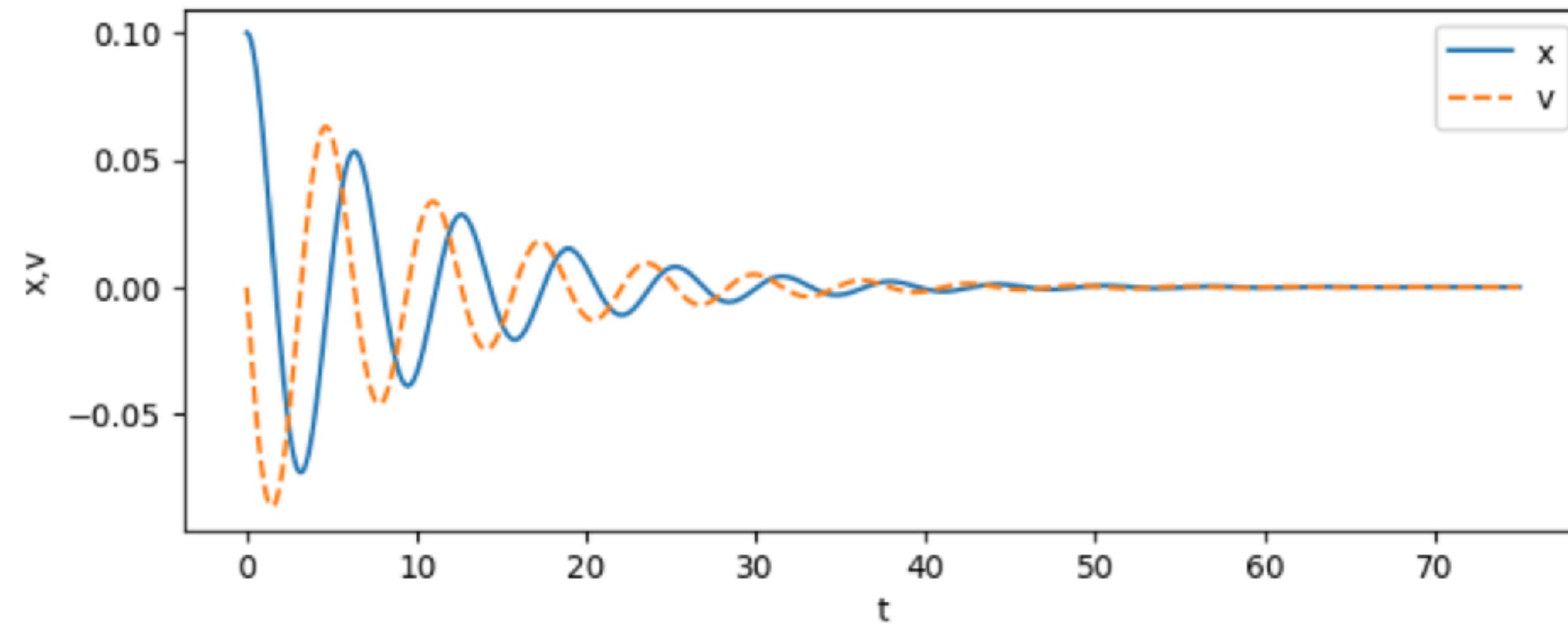
# Dynamical Systems

Example: Van der Pol oscillator:

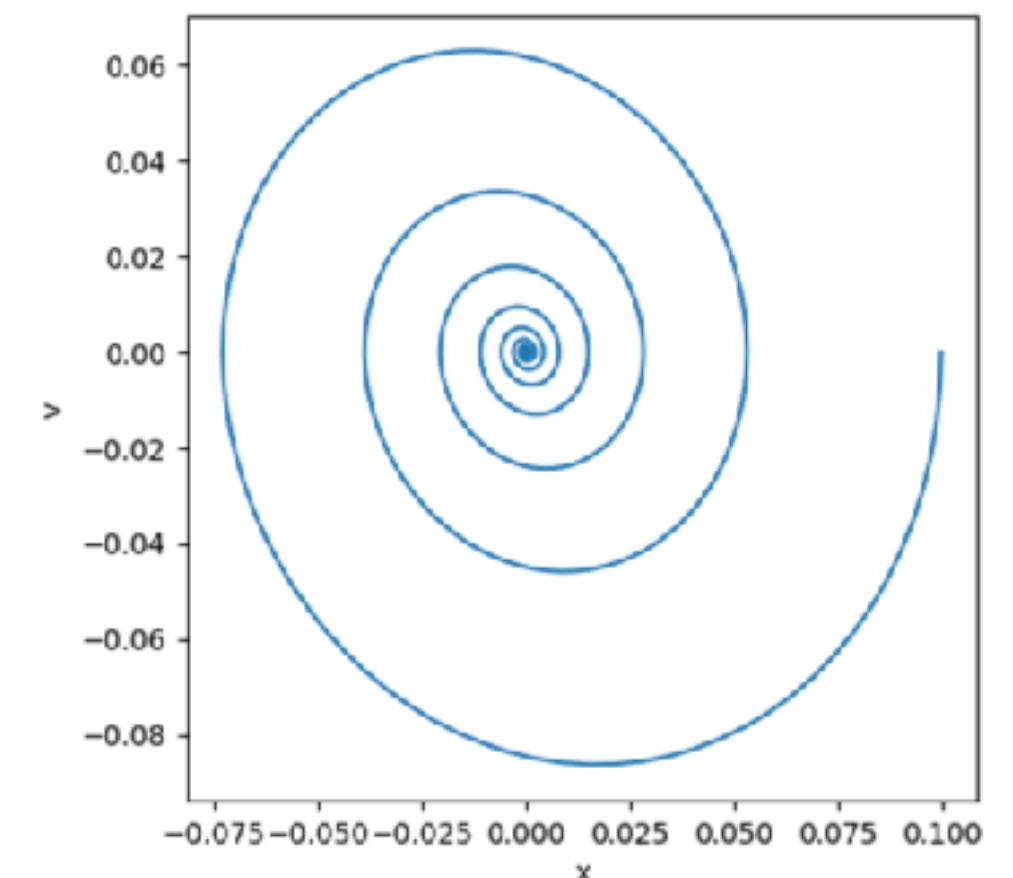
$$\frac{d^2x}{dt^2} - \mu(1 - x^2)\frac{dx}{dt} + x = 0$$

Fixed Point

$$\mu = 0.8$$

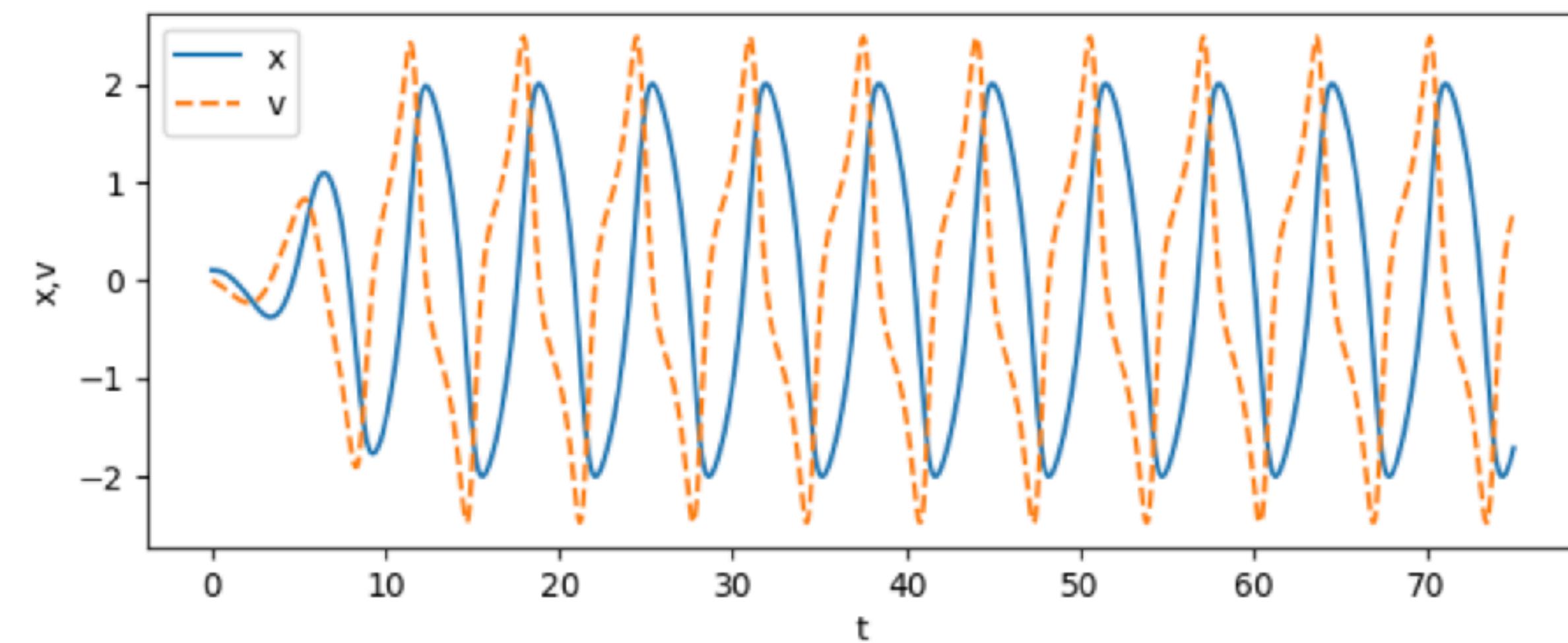


approaches single point

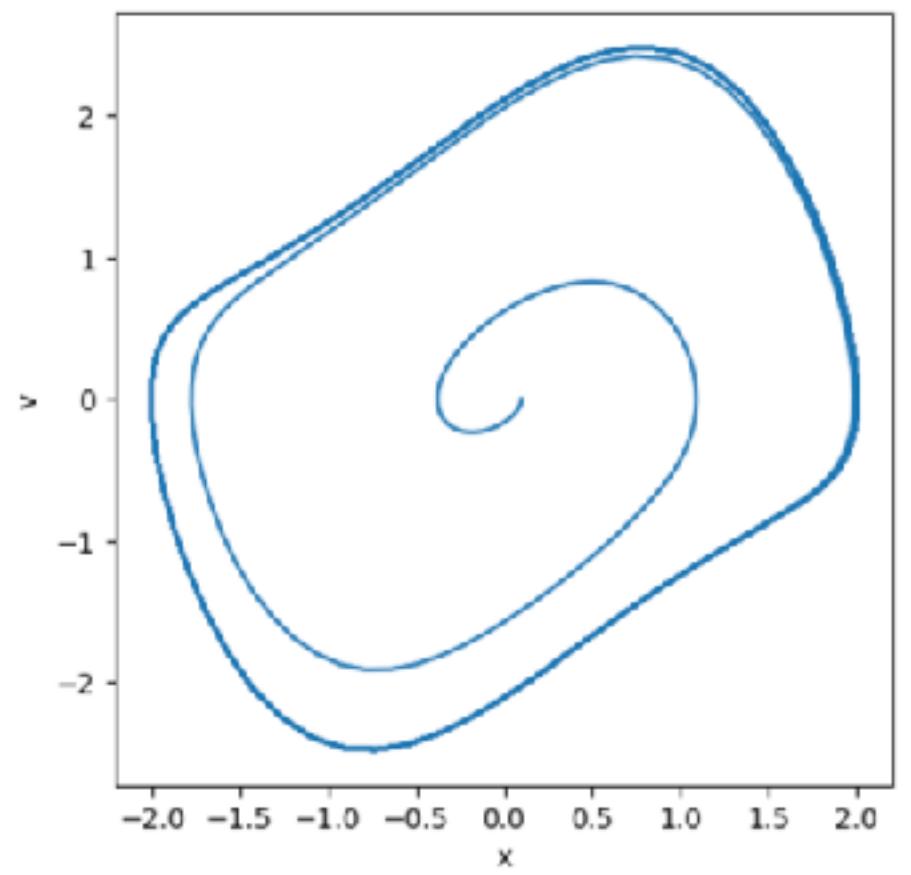


Limit Cycle

$$\mu = -0.2$$



closed orbit in phase space



# Lorenz System

Perform numerical integration to solve  
for  $x(t), y(t), z(t)$ :

- Create a derivative function
- Use the RK45 method
- Interpolate onto a uniform time array
- Standard parameter values:
  - $\sigma = 10$
  - $b = 8/3$
  - $r = 28$
- Plots:
  - time series:  $x(t), y(t), z(t)$
  - phase space:  $(x, y, z)$

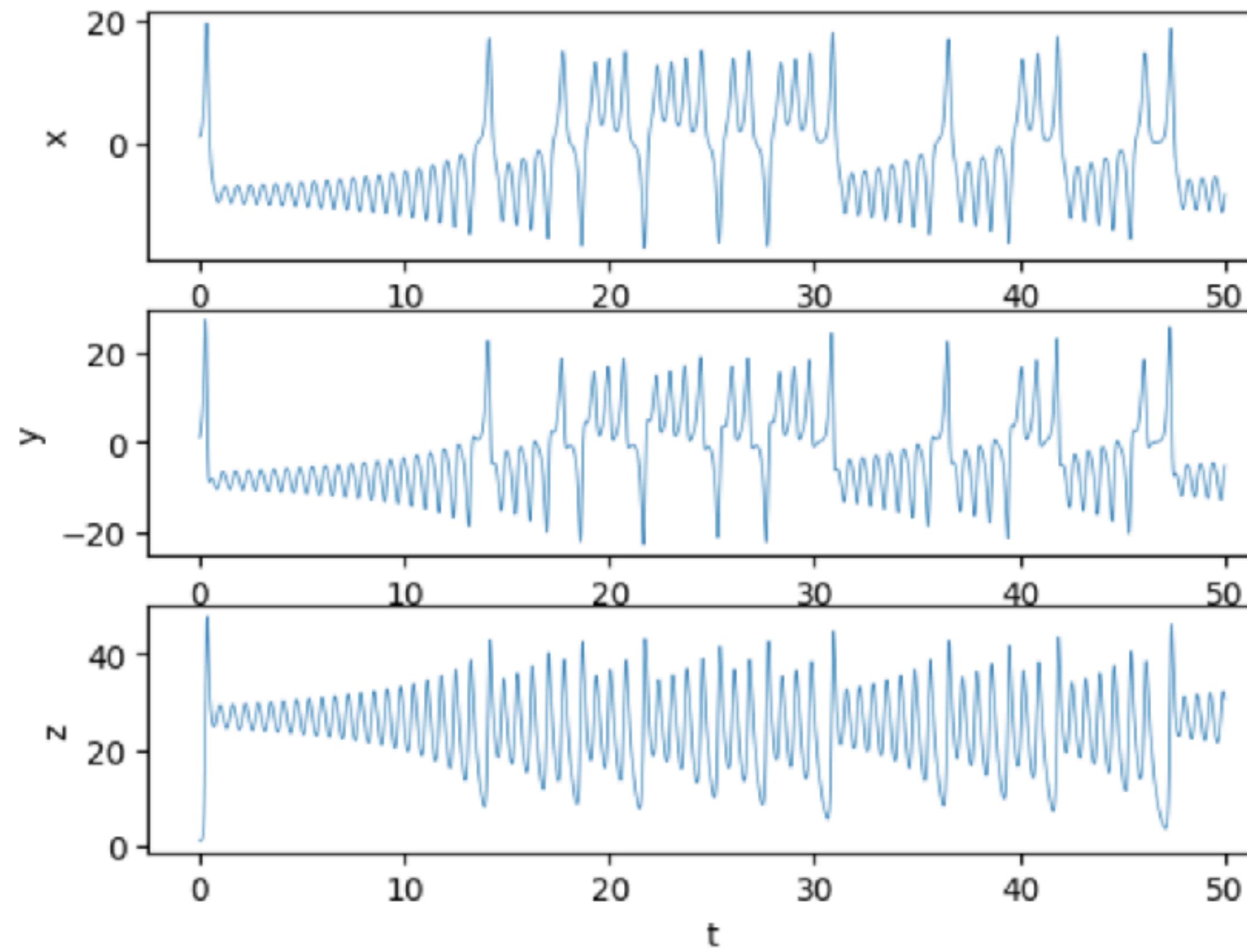
$$\frac{dx}{dt} = \sigma(y - x)$$

$$\frac{dy}{dt} = x(r - z) - y$$

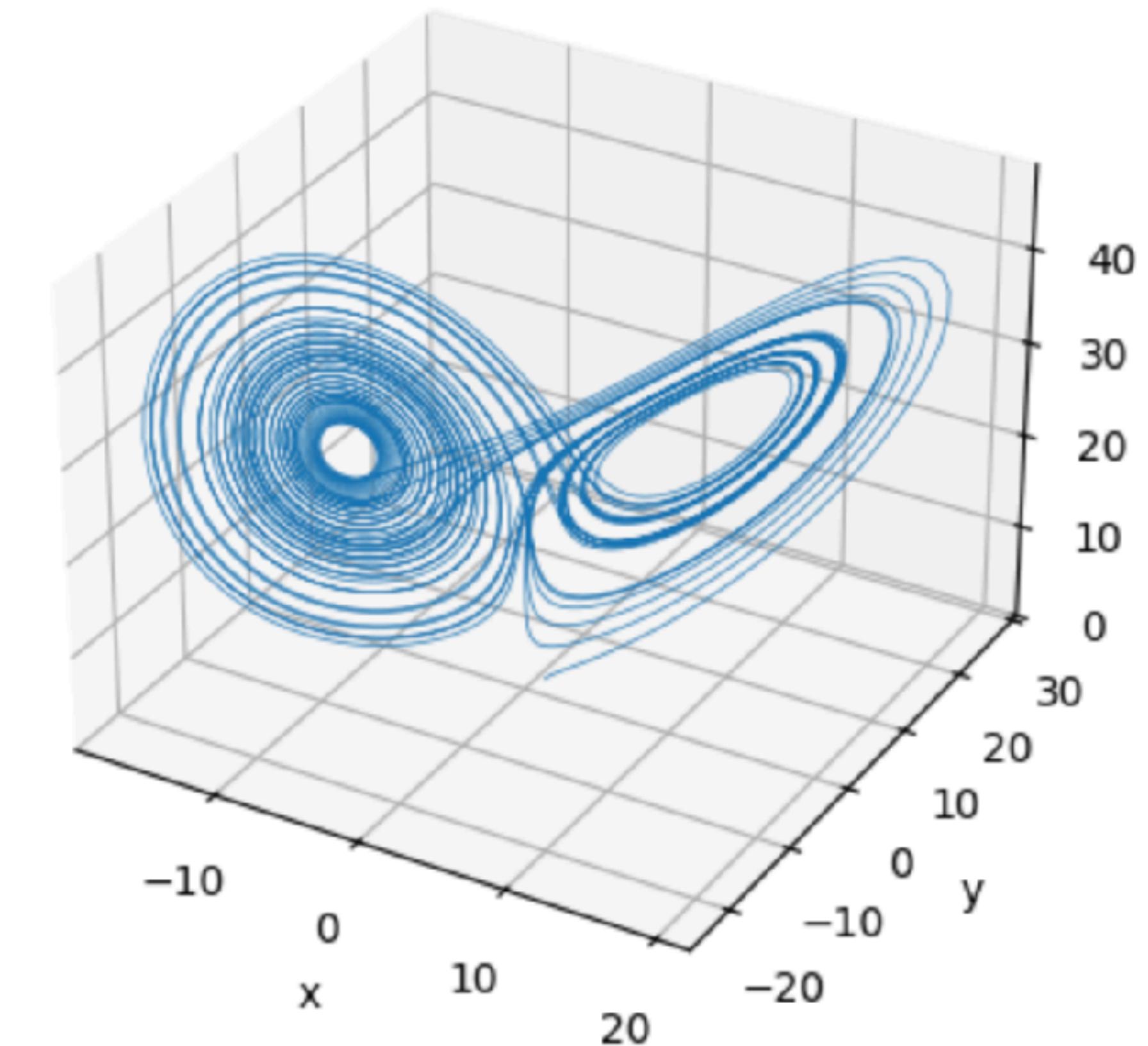
$$\frac{dz}{dt} = xy - bz$$

# Lorenz System

Lorenz System  $r = 28.00$



Dynamics do not repeat (aperiodic)

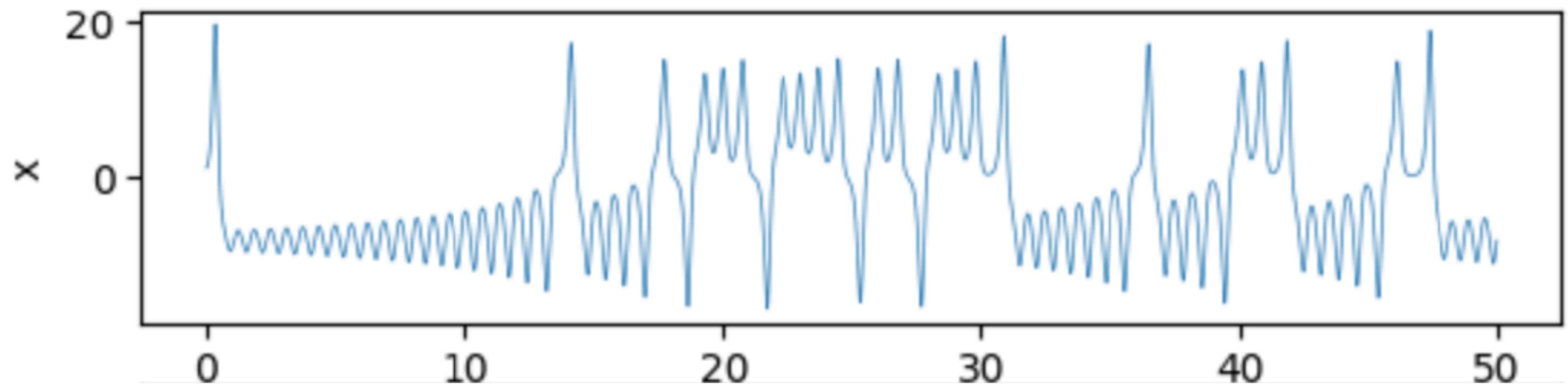


Strange attractor is neither a fixed point nor a limit cycle. It has fractal dimension

# Sensitivity to Initial Conditions

Small changes to the initial conditions can lead to large changes later

$$(x_0, y_0, z_0) = (1, 1, 1)$$



$$(x_0, y_0, z_0) = (1, 1, 1.000001)$$

$$\Delta r \propto e^{\lambda t}$$

